

## Lecture 2

### Linearity, Isotropy and Wave Solutions

## Maxwell's Equations

- $\text{Div}[\epsilon_0 \mathbf{E}] = -\text{Div}[\mathbf{P}] + \rho$
- $\text{Div}[\mu_0 \mathbf{H}] = -\text{Div}[\mu_0 \mathbf{M}]$
- $\text{Curl}[\mathbf{E}] = -\partial/\partial t (\mu_0 \mathbf{H}) - \partial/\partial t (\mu_0 \mathbf{M})$
- $\text{Curl}[\mathbf{H}] = \partial/\partial t (\epsilon_0 \mathbf{E}) + \partial/\partial t (\mathbf{P}) + \mathbf{J}$
- $\mu_0 = 4\pi \cdot 10^{-7}$  in SI units
- $\epsilon_0 = 10^{-9}/(36\pi)$  in SI units

### In The Physicist's World

- Everything is linear!!
- (In the real world everything is non-linear with very few exceptions)
- Polarisation ( $\mathbf{P}$ ), Magnetisation ( $\mathbf{M}$ ) and current density ( $\mathbf{J}$ ) are effects, not causes
- Therefore  $\mathbf{P}$  and  $\mathbf{J}$  are "proportional" to  $\mathbf{E}$
- And  $\mathbf{M}$  is "proportional" to  $\mathbf{H}$

### Tensors!!!

- Don't worry they don't last long!!
- The electric susceptibility is in general a tensor
- $\mathbf{P} = \epsilon_0 \chi_e \mathbf{E}$

$$\begin{pmatrix} \chi_{11} & \chi_{12} & \chi_{13} \\ \chi_{21} & \chi_{22} & \chi_{23} \\ \chi_{31} & \chi_{32} & \chi_{33} \end{pmatrix}$$

## Symmetry and Isotropy

- If material has no preferred axis, then it said to be **isotropic**
- If it is isotropic then the off-diagonal elements of the matrix are zero and the diagonal elements are equal
- $\mathbf{P} = \epsilon_0 \chi_e \mathbf{E}$

$$\begin{pmatrix} \chi & 0 & 0 \\ 0 & \chi & 0 \\ 0 & 0 & \chi \end{pmatrix}$$

## Same For Magnetics

- $\mathbf{M} = \chi_m \mathbf{H}$
- For isotropic media the magnetic susceptibility reduces from a tensor to a scalar
- $\mathbf{M} = \chi_m \mathbf{H}$

## Same For Current Density

- $\mathbf{J} = \sigma \mathbf{E}$
- For a linear, isotropic medium  $\sigma$  is a scalar

## Recap

- For **linear** media the response is proportional to the stimulus and the constant of proportionality is a **tensor**
- For **linear, isotropic** media the response is proportional to the stimulus and the constant of proportionality is a **scalar**
- Almost all cases are **linear**
- A large number are **linear and isotropic**

## MEs In a Linear, Isotropic Medium

- $\text{Div}[\epsilon_0 \mathbf{E}] = -\text{Div}[\chi_e \epsilon_0 \mathbf{E}] + \rho$
- $\text{Div}[\mu_0 \mathbf{H}] = -\text{Div}[\chi_m \mu_0 \mathbf{H}]$
- $\text{Curl}[\mathbf{E}] = -\partial/\partial t (\mu_0 \mathbf{H}) - \partial/\partial t (\chi_m \mu_0 \mathbf{H})$
- $\text{Curl}[\mathbf{H}] = \partial/\partial t (\epsilon_0 \mathbf{E}) + \partial/\partial t (\chi_e \epsilon_0 \mathbf{E}) + \sigma \mathbf{E}$
- This might have some solutions!!

## MEs In a Linear, Isotropic Medium

- $\text{Div}[\mathbf{E}] = \rho/\epsilon$
- $\text{Div}[\mathbf{H}] = 0$
- $\text{Curl}[\mathbf{E}] = -\partial/\partial t (\mu \mathbf{H})$
- $\text{Curl}[\mathbf{H}] = \partial/\partial t (\epsilon \mathbf{E}) + \sigma \mathbf{E}$
- $\epsilon = \epsilon_0 (1 + \chi_e)$
- $\mu = \mu_0 (1 + \chi_m)$

## MEs In a Linear, Isotropic Medium

- $\text{Curl}[\text{Curl}[\mathbf{E}]] + \partial^2/\partial t^2 (\mu \epsilon \mathbf{E}) + \partial/\partial t (\mu \sigma \mathbf{E}) = 0$
- $\text{Curl}[\text{Curl}[\mathbf{H}]] + \partial^2/\partial t^2 (\mu \epsilon \mathbf{H}) + \partial/\partial t (\mu \sigma \mathbf{H}) = 0$
- It is well-known that (ie I found it in a book somewhere)
- $\text{Curl}[\text{Curl}[\mathbf{X}]] = \text{Grad}[\text{Div}[\mathbf{X}]] - \text{Del}2[\mathbf{X}]$
- but...
- $\text{Div}[\mathbf{H}] = 0$
- so....
- $\text{Del}2[\mathbf{H}] - \partial^2/\partial t^2 (\mu \epsilon \mathbf{H}) - \partial/\partial t (\mu \sigma \mathbf{H}) = 0$
- A WAVE EQUATION!!!
- and by similar means
- $\text{Del}2[\mathbf{E}] - \partial^2/\partial t^2 (\mu \epsilon \mathbf{E}) - \partial/\partial t (\mu \sigma \mathbf{E}) = \text{Grad}[\rho/\epsilon]$
- which if  $\rho = 0$  (charge-free medium) is the same as the equation for  $\mathbf{H}$

## MEs In a Linear, Isotropic Medium

- $\text{Del}2[\mathbf{H}] - \partial^2/\partial t^2 (\mu \epsilon \mathbf{H}) - \partial/\partial t (\mu \sigma \mathbf{H}) = 0$
- $\text{Del}2[\mathbf{E}] - \partial^2/\partial t^2 (\mu \epsilon \mathbf{E}) - \partial/\partial t (\mu \sigma \mathbf{E}) = 0$
- For a **linear, isotropic, charge-free** medium
- Doesn't have to be **charge-free**, but the derivation is easier that way!!
- It doesn't even have to be **isotropic**, but the derivation is easier that way!!
- Phase speed in vacuum ( $\chi_m, \chi_e = 0$ ) is
- $(\mu_0 \epsilon_0)^{-1/2} = 3 \times 10^8 \text{ ms}^{-1} = c$

## Dissipation in MEs

- For a **linear, isotropic, charge-free** medium
- $\text{Del2}[\mathbf{H}] - \partial^2/\partial t^2 (\mu\epsilon\mathbf{H}) - \partial/\partial t (\mu\sigma\mathbf{H}) = 0$
- $\text{Del2}[\mathbf{E}] - \partial^2/\partial t^2 (\mu\epsilon\mathbf{E}) - \partial/\partial t (\mu\sigma\mathbf{E}) = 0$
- The third term in the above equation represents the dissipation
- For a **linear, isotropic, charge-free, non-conducting** medium ( $\sigma = 0$ ) the wave does not dissipate
- A Vacuum is a particular case of the above!!

## Transverse Waves?!!

- $\text{Del2}[\mathbf{H}] - \partial^2/\partial t^2 (\mu\epsilon\mathbf{H}) = 0$
- $\text{Del2}[\mathbf{E}] - \partial^2/\partial t^2 (\mu\epsilon\mathbf{E}) = 0$
- The most general plane wave solutions of the above are
- $\mathbf{E} = \mathbf{x} [ f_+(z - vt) + f_-(z + vt) ]$
- $\mathbf{H} = (\epsilon/\mu)^{1/2} \mathbf{y} [ f_+(z - vt) + f_-(z + vt) ]$
- Note that  $\mathbf{E}$  and  $\mathbf{H}$  are perpendicular to each other and to the direction of propagation - z
- But in a vacuum, these are transverse waves in what?
- And thereby hangs an aethereal tale.....

## Monochromatic Waves

- $\mathbf{E} = E_0 \mathbf{x} \cos(\mathbf{k} \cdot \mathbf{r} \pm \omega t + \phi)$
- $\phi$  is an arbitrary phase
- $\omega$  is the angular frequency
- $\mathbf{k}$  defines both the direction of propagation (DOP) and the wave vector  $|\mathbf{k}| = 2\pi/\lambda = k$
- $\mathbf{k} \cdot \mathbf{r} \pm \omega t$  defines a surface of constant phase which is displaced in time at a speed  $\omega/k$  which is called the "phase speed"

## Velocity In a Medium

- If  $\chi_e, \chi_m$  are  $\neq 0$ , then the phase speed is reduced
- Phase speed becomes  $c/n$
- $n = (\epsilon/\epsilon_0 \cdot \mu/\mu_0)^{1/2}$
- $n$  is the "refractive index" and is  $> 1$  (mostly)
- If either of  $\chi_e, \chi_m$  are dependent upon frequency, then the medium has "dispersion"
- This means that if  $n$  is frequency-dependent, then the medium is dispersive

## Complex Numbers

- Complex numbers for physicists are a “convenience” for solving problems
- Rely on the real-world fiction that  $\sqrt{-1}$  can be represented by  $i$  (the engineers use  $j$ )
- A complex number can be written as  $C = a + ib$
- The conjugate complex is  $C^* = a - ib$
- The modulus is  $|C| = \sqrt{a^2 + b^2}$
- The argument is  $\text{Arg}(C) = \tan^{-1}(b/a)$

## Complex Exponentials - Phasors

- If  $C = \cos(a) + i \sin(a)$
- Then it is shown in all the best text books that you can write  $C = \exp(ia)$
- In fact any complex number  $C$  can be written as  $C = |C| \exp(i \arg(C))$
- Since manipulation of exponentials in integration and differentiation is very easy, it is tempting to try to manipulate any wave equation into this format

## Complex Superposition

- If  $E_0 \mathbf{x} \cos(\mathbf{k} \cdot \mathbf{r} - \omega t + \phi)$  is a solution
- Then so is  $E_0 \mathbf{x} \sin(\mathbf{k} \cdot \mathbf{r} - \omega t + \phi)$
- And so is  $i \cdot E_0 \mathbf{x} \sin(\mathbf{k} \cdot \mathbf{r} - \omega t + \phi)$
- By the principle of linear superposition, the sum of the above solutions is also a solution which can be expressed as
- $E_0 \mathbf{x} \exp i(\mathbf{k} \cdot \mathbf{r} - \omega t + \phi)$
- or  $\mathbf{x} \cdot E \cdot \exp(-i\omega t)$  where  $E$  is complex and contains all spatial and phase components  $E = E_0 \exp i(\mathbf{k} \cdot \mathbf{r} + \phi)$

## Differentiation wrt Time

- For any quantity  $C \exp(-i\omega t)$  where  $C$  is independent of time
- $\partial/\partial t (C \exp(-i\omega t)) = -i\omega C \exp(-i\omega t)$
- Similarly  $\text{Int}[C \exp(-i\omega t)] = -1/i\omega C \exp(-i\omega t)$
- What does this do to Maxwell's equations?
- $\text{Div}[\epsilon_0 \mathbf{E}] = -\text{Div}[\mathbf{P}] + \rho$  - unchanged
- $\text{Div}[\mu_0 \mathbf{H}] = -\text{Div}[\mu_0 \mathbf{M}]$  - unchanged
- $\text{Curl}[\mathbf{E}] = -\partial/\partial t (\mu_0 \mathbf{H}) - \partial/\partial t (\mu_0 \mathbf{M})$  becomes
- $\text{Curl}[\mathbf{E}] = i\omega\mu_0 \mathbf{H} + i\omega\mu_0 \mathbf{M}$
- $\text{Curl}[\mathbf{H}] = \partial/\partial t (\epsilon_0 \mathbf{E}) + \partial/\partial t (\mathbf{P}) + \mathbf{J}$  becomes
- $\text{Curl}[\mathbf{H}] = -i\omega\epsilon_0 \mathbf{E} - i\omega \mathbf{P} + \mathbf{J}$

## ME Monochromatic Wave Eqn

- $\text{Del}^2[ \mathbf{H} ] + \omega^2 \mu \epsilon \mathbf{H} + i \omega \mu \sigma \mathbf{H} = 0$
- $\text{Del}^2[ \mathbf{E} ] + \omega^2 \mu \epsilon \mathbf{E} + i \omega \mu \sigma \mathbf{E} = 0$
- $\mu_0 = 4\pi \cdot 10^{-7}$  in SI units
- $\epsilon_0 = 10^{-9}/(36\pi)$  in SI units
- If we now REDEFINE  $\epsilon$  as being  $\epsilon + i\sigma/\omega$ , we can automatically incorporate conducting media in the solutions with a complex dielectric constant
- This implies that the effect of conductance is the the polarisation  $\mathbf{P}$  is still proportional to the field  $\mathbf{E}$ , but there is a phase lag between the field and the polarisation
- We can also look at the same arguments for the spatial differential  $\text{Del}^2[ ]$  and find that in a charge-free medium  $\text{Del}^2[\mathbf{X}] = -k^2 \mathbf{X}$  where  $k$  is the wave vector

## Dispersion Relation

- Putting all these components into the wave equation, we find that
- $k^2 = \omega^2 \epsilon \mu$  for a linear, isotropic, charge-free medium
- Remember that  $\epsilon$  is generally complex and therefore so is  $k$
- Remember also that  $\mu \approx \mu_0$  in most cases
- If we define refractive index to be complex following  $\epsilon$ , then
- $k^2 = \omega^2 / c^2 \cdot n^2$
- This is the “dispersion relationship” and measures how the wave vector  $k$ , varies with frequency

## n and $\epsilon$

- Most materials are non-magnetic, assume  $\mu = \mu_0$
- $\epsilon = \epsilon_r + i \epsilon_i = n^2 \epsilon_0$
- So if  $n = n_r + i n_i$
- Then  $\epsilon_r / \epsilon_0 = n_r^2 - n_i^2$
- And  $\epsilon_i / \epsilon_0 = 2 n_r n_i$
- What's the point?
- We can measure the refractive index very easily
- therefore we can relate the macroscopic  $n$  to the microscopic, complex  $\epsilon$

## Bringing It Back to Reality

- The wave propagates like  $\exp i(\mathbf{k} \cdot \mathbf{r} - \omega t + \phi)$
- Suppose it propagates along the  $z$  axis  $\rightarrow kz$
- The wavelength in free space is  $\lambda_v$
- and the wavevector is  $k_v$
- Look at the spatial part only
- $\exp(ikz) = \exp\{i (n_r + i n_i) k_v z\}$
- $\exp(ikz) = \exp(i n_r k_v z) \exp(-n_i k_v z)$
- First term is phase factor, second is a decay term
- Measure the phase speed - gives  $n_r$  (historically  $n$ )
- Measure the decay - gives  $n_i$  (historically  $\kappa$ )

## Refractive Indexes

- Vacuum = 1.00 (surprise, surprise)
- Glasses = 1.57-1.77
- Sapphire = 1.77, Diamond = 2.42
- Silicon =  $3.8 + 0.4i$
- Silver =  $2.3 + 3.8i$
- Air =  $1.0 + 290 \cdot 10^{-6} \left( \rho/\rho_{\text{stp}} \right)$